

# ASYMPTOTICS OF $10j$ SYMBOLS

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ABSTRACT. The Riemannian  $10j$  symbols are spin networks that assign an amplitude to each 4-simplex in the Barrett-Crane model of Riemannian quantum gravity. This amplitude is a function of the areas of the 10 faces of the 4-simplex, and Barrett and Williams have shown that one contribution to its asymptotics comes from the Regge action for all non-degenerate 4-simplices with the specified face areas. However, we show numerically that the dominant contribution comes from degenerate 4-simplices. As a consequence, one can compute the asymptotics of the Riemannian  $10j$  symbols by evaluating a ‘degenerate spin network’, where the rotation group  $SO(4)$  is replaced by the Euclidean group of isometries of  $\mathbb{R}^3$ . We conjecture formulas for the asymptotics of a large class of Riemannian and Lorentzian spin networks in terms of these degenerate spin networks, and check these formulas in some special cases. Among other things, this conjecture implies that the Lorentzian  $10j$  symbols are asymptotic to  $1/16$  times the Riemannian ones.

## 1. INTRODUCTION

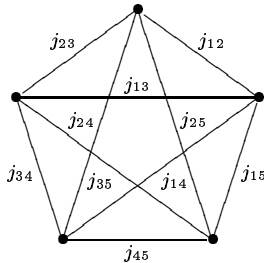
In the Ponzano–Regge model of 3-dimensional Riemannian quantum gravity [1], an amplitude is associated with each tetrahedron in a triangulation of spacetime. The amplitude depends on the tetrahedron’s six edge lengths, which are assumed to be quantized, taking values proportional to  $2j + 1$  where  $j$  is a half-integer spin. One can compute this amplitude either by evaluating an  $SU(2)$  spin network shaped like a tetrahedron, or by doing an integral. Approximating this integral by the stationary phase method, Ponzano and Regge argued that when all six spins are rescaled by the same factor  $\lambda$ , the  $\lambda \rightarrow \infty$  asymptotics of the amplitude are given by a simple function of the volume of the tetrahedron and the Regge calculus version of its Einstein action. Nobody has yet succeeded in making their argument rigorous, but Roberts [2, 3] recently proved their asymptotic formula by a different method. This result lays the foundation for a careful study of the relation between the Ponzano–Regge model and classical general relativity in 3 dimensions.

Our concern here is whether a similar result holds for the Barrett–Crane model of *4-dimensional* Riemannian quantum gravity [4]. In this model an amplitude is associated with each 4-simplex in a triangulation of spacetime. This amplitude, known as a  $10j$  symbol, is a function of the areas of the 10 triangular faces of the 4-simplex. Each triangle area is proportional to  $2j + 1$ , where  $j$  is a spin labelling the triangle. The amplitude can be computed by evaluating an  $SU(2) \times SU(2)$  spin network whose edges correspond to the

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triangles of the 4-simplex:



There is also an integral formula for the  $10j$  symbol [7]. The problem is to understand the asymptotics of the  $10j$  symbol as all 10 spins are rescaled by a factor  $\lambda$  and  $\lambda \rightarrow \infty$ .

Barrett and Williams [8] applied a stationary phase approximation to the integral for the  $10j$  symbol, focusing attention on stationary phase points corresponding to nondegenerate 4-simplices with the specified face areas. They showed that each such 4-simplex contributes to the  $10j$  symbol in a manner that depends on its Regge action. They pointed out the existence of contributions from degenerate 4-simplices, but did not analyse them.

In Section 2 of this paper we begin by applying Barrett and Williams' estimate of the  $10j$  symbols to the case where all 10 spins are equal. We find that the contribution of their stationary phase points to the  $10j$  symbols is of order  $\lambda^{-9/2}$ . However, our numerical calculations show that the  $10j$  symbols are much larger, of order  $\lambda^{-2}$ . This means we must look elsewhere to explain the asymptotics of the  $10j$  symbol.

In Section 3 we analyse the contribution of 'degenerate 4-simplices' to the integral for the  $10j$  symbols. They do not correspond to stationary phase points in the integral for the  $10j$  symbols; instead, the integrand has a strong *maximum* at these points. We argue that the contribution of a small neighborhood of these points is asymptotically proportional to  $\lambda^{-2}$ . We give a formula expressing the constant of proportionality as an integral over the space of degenerate 4-simplices. We also reduce this to an explicit integral in 5 variables.

In Section 4, we numerically compare these results to the  $10j$  symbols as calculated using the algorithm developed by Christensen and Egan [9]. Our formula for the contribution of degenerate 4-simplices closely matches the actual asymptotics of the  $10j$  symbols. Thus, even though our argument that these asymptotics are dominated by degenerate 4-simplices is not rigorous, we feel confident that the resulting formula is correct.

In Section 5 we discuss a new sort of spin network, associated to the representation theory of the Euclidean group, which arises naturally in our analysis of the contribution of degenerate 4-simplices. Generalizing our results on the Riemannian  $10j$  symbols, we conjecture formulas for the asymptotics of a large class of Riemannian spin networks in terms of these new 'degenerate spin networks'. We verify this conjecture in a number of simple cases.

In Section 6 we formulate a similar conjecture for Lorentzian spin networks. Taken with the previous one this conjecture implies that the  $\lambda \rightarrow \infty$  asymptotics of a Lorentzian spin network in this class are the same, up to a constant, as those of the corresponding Riemannian spin network. For example, as  $\lambda \rightarrow \infty$ , the Lorentzian  $10j$  symbol should be asymptotic to  $1/16$  times the corresponding Riemannian  $10j$  symbol! We conclude by presenting some numerical evidence that this is the case.

## 2. STATIONARY PHASE POINTS

The  $10j$  symbols can be defined using a Riemannian spin network — also known as a 'balanced' spin network [4] — whose underlying graph is the complete graph on five vertices.

The ten edges of the graph are labelled with half-integer spins  $j_{kl} = 0, \frac{1}{2}, 1, \dots$ , where  $k$  and  $l$  refer to the vertices connected by each edge. In this approach, an edge labelled by the spin  $j$  corresponds to the representation  $j \otimes j$  of  $\text{Spin}(4) = \text{SU}(2) \times \text{SU}(2)$ , and the  $10j$  symbols are computed using the representation theory of this group.

However, to analyse the asymptotics of the  $10j$  symbol, it is easier to use the integral formula due to Barrett [7]. This is:

$$(1) \quad \begin{array}{c} \begin{array}{c} \text{Diagram of a 5-vertex graph with 10 edges labeled } j_{23}, j_{12}, j_{13}, j_{24}, j_{25}, j_{34}, j_{35}, j_{14}, j_{15}, j_{45} \end{array} \\ \end{array} = (-1)^{\sum_{k<l} 2j_{kl}} \int_{(S^3)^5} \prod_{k<l} K_{2j_{kl}+1}^R(\phi_{kl}) \frac{dh_1}{2\pi^2} \cdots \frac{dh_5}{2\pi^2},$$

where  $S^3$  is the unit 3-sphere in  $\mathbb{R}^4$  equipped with its usual Lebesgue measure,  $\phi_{kl}$  is the angle between the unit vectors  $h_k$  and  $h_l$ , and the kernel  $K^R$  is given by:

$$(2) \quad K_a^R(\phi) := \frac{\sin a\phi}{\sin \phi}.$$

The normalizing factors in the integral come from the fact that the volume of the unit 3-sphere is  $2\pi^2$ .

The  $10j$  symbol gives the amplitude for a 4-simplex with specified triangle areas. Each vertex of the above graph corresponds to a tetrahedron in this 4-simplex, and each edge of the graph corresponds to the unique triangle shared by two of these tetrahedra. In this picture, the spin  $j_{kl}$  determines the area of the triangle shared by the  $k$ th and  $l$ th tetrahedra. The precise formula for this area is somewhat controversial [10, 11], but given the integral formula for the  $10j$  symbols, we find it convenient to assume the area is proportional to  $2j_{kl} + 1$ . Ignoring the constant factor, we thus define triangle areas by:

$$a_{kl} = 2j_{kl} + 1.$$

In what follows, we study the behaviour of the  $10j$  symbol as all these triangle areas  $a_{kl}$  are multiplied by a large integer  $\lambda$ . This is not the same as multiplying the spins  $j_{kl}$  by  $\lambda$ . However, note that as  $j_{kl}$  ranges over all spins,  $a_{kl}$  ranges over all positive integers. This means that if we multiply the areas  $a_{kl}$  by any positive integer  $\lambda$ , we can find new spins  $J_{kl}$  corresponding to the new areas by solving  $\lambda a_{kl} = 2J_{kl} + 1$ .

It is shown in [5] that the integral in (1) is nonnegative. Accordingly, we will concentrate on analysing the asymptotics of the absolute value of the  $10j$  symbol, given by the integral alone:

$$(3) \quad \left| \begin{array}{c} \text{Diagram of a 5-vertex graph with 10 edges} \\ \end{array} \right| := \int_{(S^3)^5} \prod_{k<l} K_{\lambda a_{kl}}^R(\phi_{kl}) \frac{dh_1}{2\pi^2} \cdots \frac{dh_5}{2\pi^2},$$

where for simplicity we have left out the spins labelling the spin network edges, which are now  $J_{kl}$ .

Barrett and Williams [8] express the numerator in the kernel  $K^R$  as a difference of exponentials, which allows them to rewrite the integrand in equation (3) as a sum of  $2^{10}$  terms, each consisting of an exponential times the function:

$$(4) \quad f(h_1, \dots, h_5) = \prod_{k<l} \frac{1}{\sin \phi_{kl}}.$$

This function is unbounded as any of the  $\phi_{kl}$  tend to zero or  $\pi$ , but if the regions of the domain where that occurs are set aside for a separate analysis, the integral over the remaining region becomes amenable to a stationary phase approximation [14]. The relevant phase is simply:

$$(5) \quad S(h_1, \dots, h_5) = \sum_{k < l} \lambda a_{kl} \phi_{kl}.$$

The phases of the  $2^{10}$  exponentials arising from the product of the kernels include variants of  $S$  with all possible signs for the ten terms. However, Barrett and Williams use the invariance of the integrand under the transformations  $h_k \rightarrow -h_k$  to break the integral into  $2^5$  identical parts, according to whether  $h_k = \pm n_k$ , where  $n_k$  are the outward normals of a 4-simplex whose five tetrahedra lie in the hyperplanes normal to the  $h_k$ . They then focus on the region where all the  $h_k = n_k$ , and show that of the  $2^{10}$  phases, only  $S$  and its opposite have stationary points in this region, and they occur when the  $\phi_{kl}$  are the angles between the outward normals to the tetrahedra of a 4-simplex whose ten faces have areas given by  $\lambda a_{kl}$ . In this case the  $h_k$  can be interpreted as the outward normals to the five tetrahedra, the  $\phi_{kl}$  are the angles between these normals, and  $\lambda a_{kl}$  is the area of the face shared by the tetrahedra numbered  $k$  and  $l$ . Interpreted this way,  $S$  is precisely the Regge action for the 4-simplex.

This means that the relevant stationary phase points are those of a much simpler integral:

$$(6) \quad \frac{1}{2^5} \int f(h_1, \dots, h_5) (e^{iS(h_1, \dots, h_5)} + e^{-iS(h_1, \dots, h_5)}) \frac{dh_1}{2\pi^2} \dots \frac{dh_5}{2\pi^2}.$$

To carry out a stationary phase approximation of this integral we must note that the stationary phase ‘point’ determined by the geometry of a 4-simplex is not actually a single point in the 15-dimensional manifold  $(S^3)^5$ , but rather a whole 6-dimensional submanifold, since the geometry of the 4-simplex is invariant under the 6-parameter rotation group  $\text{SO}(4)$ . However, the same invariance can be used to remove this complication. Since the functions  $f$  and  $S$  are invariant under the action of  $\text{SO}(4)$ , they pass to the quotient space  $(S^3)^5 / \text{SO}(4)$ , and we obtain:

$$(7) \quad \frac{1}{2(2\pi)^{10}} \int_{(S^3)^5 / \text{SO}(4)} f(x) (e^{iS(x)} + e^{-iS(x)}) d\mu(x),$$

where  $d\mu$  is the result of pushing forward Lebesgue measure on  $(S^3)^5$  to this quotient space. A factor of  $1/2$  here accounts for the complication that simultaneously negating all the  $h_k$  has the same action as  $-I$ , an element of  $\text{SO}(4)$ , and so the discrete symmetries used to obtain (6) already entailed taking the quotient by a 2-element subgroup of  $\text{SO}(4)$ .

Now, the stationary phase approximation [14] of an  $n$ -dimensional integral

$$\int f(x) e^{iS(x)} dx$$

is a sum over stationary points  $x_i$  of the function  $S$ :

$$(2\pi)^{n/2} \sum_i \frac{f(x_i)}{|\det H(x_i)|^{1/2}} \exp \left[ iS(x_i) + \frac{i\pi}{4} \sigma(H(x_i)) \right],$$

where  $H(x_i)$  is the matrix of second partial derivatives of  $S$  at the point  $x_i$ , and  $\sigma(H(x_i))$  is the signature of this matrix, i.e., the number of positive eigenvalues minus the number of negative eigenvalues. This approximation assumes there are finitely many stationary points, all with  $\det H(x_i) \neq 0$ . Applying this to the case at hand, and neglecting points where some

of the angles  $\phi_{kl}$  are 0 or  $\pi$ , we obtain Barrett and Williams' stationary phase approximation of the  $10j$  symbols:

$$(8) \quad \left| \begin{array}{c} \text{Diagram of a 4-simplex with 5 vertices and 10 edges} \\ \text{stat} \end{array} \right| := \frac{1}{(2\pi)^{11/2}} \sum_i \frac{f(x_i)}{|\det H(x_i)|^{1/2}} \cos \left[ S(x_i) + \frac{\pi}{4} \sigma(H(x_i)) \right],$$

where the sum is over the 4-simplices  $x_i$  with the specified face areas.

We can easily say something about the  $\lambda \rightarrow \infty$  behavior of this quantity without actually evaluating it. Each stationary point  $x_i$  is independent of  $\lambda$ , so  $f(x_i)$  will be constant as  $\lambda \rightarrow \infty$ , while  $S(x_i)$  will grow linearly, as will its matrix of second derivatives,  $H(x_i)$ . This is a  $9 \times 9$  matrix, since  $(S^3)^5/\text{SO}(4)$  is 9-dimensional, so its determinant will grow as  $\lambda^9$ . It follows that:

$$(9) \quad \left| \begin{array}{c} \text{Diagram of a 4-simplex with 5 vertices and 10 edges} \\ \text{stat} \end{array} \right| = O(\lambda^{-9/2}).$$

However, cancellation between different stationary points could in principle make this a misleading overestimate. Thus it seems worthwhile to explicitly evaluate the sum in equation (8), at least in a simple special case.

In what follows we evaluate this sum for the case of a  $10j$  symbol with all edges labelled by the same spins. In other words, we calculate the stationary phase approximation of the amplitude for a regular 4-simplex.

The first step, which would be useful more generally, is to find explicit coordinates on the quotient space  $(S^3)^5/\text{SO}(4)$  and describe the measure  $d\mu$  in these terms. To do this, we exploit the fact that any configuration of the five unit vectors  $h_k$  can be rotated into one that belongs to a 9-dimensional subspace of the original domain. Specifically, if we express the vectors  $h_k$  in polar coordinates

$$h_k = (\cos \psi_k, \sin \psi_k \cos \theta_k, \sin \psi_k \sin \theta_k \cos \phi_k, \sin \psi_k \sin \theta_k \sin \phi_k),$$

where  $0 \leq \psi_k, \theta_k \leq \pi$  and  $0 \leq \phi_k \leq 2\pi$ , any set of  $h_k$  can be rotated in such a way that the following restrictions are met:

$$(10a) \quad h_1 = (1, 0, 0, 0) \quad \psi_1 = \theta_1 = \phi_1 = 0$$

$$(10b) \quad h_2 = (\cos \psi_2, \sin \psi_2, 0, 0) \quad \theta_2 = \phi_2 = 0$$

$$(10c) \quad h_3 = (\cos \psi_3, \sin \psi_3 \cos \theta_3, \sin \psi_3 \sin \theta_3, 0) \quad \phi_3 = 0.$$

By means of this 'gauge-fixing' we can use the remaining 9 variables  $\psi_2, \psi_3, \psi_4, \psi_5, \theta_3, \theta_4, \theta_5, \phi_4, \phi_5$  as coordinates on the quotient space. To describe the measure  $d\mu$  in these coordinates recall that in polar coordinates, Lebesgue measure on the unit 3-sphere is given by

$$(11a) \quad dh_k = \sin^2 \psi_k \sin \theta_k d\psi_k d\theta_k d\phi_k.$$

Since  $h_1$  is completely fixed we omit  $dh_1$  from the formula for  $d\mu$ , only inserting a factor  $2\pi^2$  due to the volume of the 3-sphere. Since  $h_2$  and  $h_3$  are partially fixed we replace  $dh_2$  and  $dh_3$  by

$$(11b) \quad \tilde{d}h_2 = 4\pi \sin^2 \psi_2 d\psi_2$$

$$(11c) \quad \tilde{d}h_3 = 2\pi \sin^2 \psi_3 \sin \theta_3 d\psi_3 d\theta_3.$$

We thus obtain

$$(12) \quad d\mu = 2\pi^2 \tilde{d}h_2 \tilde{d}h_3 dh_4 dh_5.$$

Next, we determine the stationary points of the function  $S$  in the case where all ten spins are equal, neglecting points where some of the angles  $\phi_{kl}$  are 0 or  $\pi$  — that is, where some of the vectors  $h_k$  are parallel or anti-parallel. Recall that these stationary points correspond to nondegenerate 4-simplices having all 10 face areas equal. One obvious candidate is the regular 4-simplex. However, we must rule out the possibility that there are other, *non-regular* 4-simplices for which all the faces have identical areas. Since we are excluding degenerate 4-simplices from the current analysis, we can appeal to a theorem of Bang [15] which states that if all the faces of a non-degenerate tetrahedron have the same area, they are all congruent. It follows that if all ten triangles in a non-degenerate 4-simplex have the same area, they too are all congruent. Now, each of the ten edges of a 4-simplex is shared by three triangles, so we can treat each edge as a triple of congruent line segments that happen to be superimposed, giving a total of 30 in all. If all the triangles were congruent isosceles triangles, each with one side of length  $L$ , then ten of these 30 line segments would be of length  $L$ . However, there is no way to partition ten line segments into congruent triples. The same argument rules out scalene triangles. So the faces must be equilateral, and the 4-simplex must be regular.

We must therefore find all sets of unit vectors  $h_k$  which satisfy the gauge-fixing conditions in equation (10) and form the outward normals of a regular 4-simplex. One choice is:

$$\begin{aligned} n_1 &= (1, 0, 0, 0) \\ n_2 &= \left(-\frac{1}{4}, \frac{\sqrt{15}}{4}, 0, 0\right) \\ n_3 &= \left(-\frac{1}{4}, -\frac{\sqrt{5/3}}{4}, \sqrt{5/6}, 0\right) \\ n_4 &= \left(-\frac{1}{4}, -\frac{\sqrt{5/3}}{4}, -\frac{\sqrt{5/6}}{2}, -\frac{\sqrt{5/2}}{2}\right) \\ n_5 &= \left(-\frac{1}{4}, -\frac{\sqrt{5/3}}{4}, -\frac{\sqrt{5/6}}{2}, \frac{\sqrt{5/2}}{2}\right). \end{aligned}$$

These vectors have mutual dot products of  $-\frac{1}{4}$ . They represent a single point in our 9-dimensional domain:

$$(13a) \quad \psi_2 = \psi_3 = \psi_4 = \psi_5 = \cos^{-1}\left(-\frac{1}{4}\right)$$

$$(13b) \quad \theta_3 = \theta_4 = \theta_5 = \cos^{-1}\left(-\frac{1}{3}\right)$$

$$(13c) \quad \phi_4 = -\cos^{-1}\left(-\frac{1}{2}\right) = -\frac{2\pi}{3}$$

$$(13d) \quad \phi_5 = \cos^{-1}\left(-\frac{1}{2}\right) = \frac{2\pi}{3}.$$

The only other choice comes from interchanging  $\phi_4$  and  $\phi_5$ . This yields another regular 4-simplex.

Now we take the phase (5) and specialise to the case where all the spins  $j_{kl}$  are equal, say to  $j$ . Setting  $a = 2j + 1$ , and working in our chosen coordinates, we obtain

$$(14) \quad S(\psi_2, \psi_3, \psi_4, \psi_5, \theta_3, \theta_4, \theta_5, \phi_4, \phi_5) = \lambda a \sum_{k < l} \cos^{-1}(h_k \cdot h_l)$$

The partial derivatives of  $S$  are all zero at the point described by (13). Symbolic computer calculations show that at this point the matrix of second derivatives of  $S$  has 5 positive eigenvalues and 4 negative ones, and a determinant of:

$$\sqrt{\frac{5}{3}} \cdot \frac{5}{2} \left(\frac{2}{3}\right)^{11} (\lambda a)^9.$$

Applying this result to (8) at the point described by (13), working in our chosen coordinates, redefining the function  $f$  to include the measure (12) as well as the kernel denominators (4),

and multiplying by a factor of two to account for the second point where  $\phi_4$  and  $\phi_5$  are interchanged, we obtain:

$$(15) \quad \left| \begin{array}{c} \text{Diagram} \\ \text{stat} \end{array} \right| = \frac{24\left(\frac{3}{5}\right)^{\frac{3}{4}}}{5\pi^{\frac{3}{2}}} (\lambda a)^{-9/2} \cos \left[ 10 \cos^{-1} \left(-\frac{1}{4}\right) \lambda a + \frac{\pi}{4} \right]$$

The above expression consists of an oscillating term times a function proportional to  $\lambda^{-9/2}$ . In short, the estimate

$$\left| \begin{array}{c} \text{Diagram} \\ \text{stat} \end{array} \right| = O(\lambda^{-9/2})$$

is sharp, at least for the case of 10 equal spins.

However, our numerical calculations of the  $10j$  symbols exhibit very different behaviour. Rather than being of order  $\lambda^{-9/2}$ , they appear to be much larger, of order  $\lambda^{-2}$ . Also, they exhibit no discernable oscillations. In the next section we argue that these results are explained by the contribution of ‘degenerate 4-simplices’.

### 3. DEGENERATE POINTS

The absolute value of the kernel

$$(16) \quad K_a^R(\phi) = \frac{\sin a\phi}{\sin \phi}$$

has maxima when the angle  $\phi$  equals 0 or  $\pi$ , and these maxima become ever more sharply peaked as  $a \rightarrow \infty$ . This means that as  $\lambda \rightarrow \infty$ , the integrand in equation (3) becomes very large at ‘degenerate points’, where some of the vectors  $h_k$  are either parallel or anti-parallel. We have made a detailed study of the ‘fully degenerate’ points, where *all* the vectors  $h_k$  are either parallel or anti-parallel. It is plausible to expect the contribution from a neighborhood of these points to dominate the integral for the  $10j$  symbol, at least asymptotically as  $\lambda \rightarrow \infty$ , since in this region we are integrating a product of kernels, all of which are near their greatest possible absolute value. In the next section we shall present numerical evidence that this is in fact the case.

The value of the kernel is positive at  $\phi = 0$ , but at  $\phi = \pi$  its sign is positive when  $2j$  is even and negative otherwise, where  $a = 2j + 1$ . At first glance, this appears to allow the possibility that parallel and anti-parallel degenerate points might cancel each other for certain values of the spins. In fact, cancellation can occur only when the  $10j$  symbol vanishes. The integrand in (1) is the product of ten kernels, so it will be positive when all the  $h_k$  are parallel. If  $h_1$  is replaced by its opposite, the four  $\phi_{1l}$  will change from zero to  $\pi$ , leading to an overall sign change of  $(-1)^{2(j_{12}+j_{13}+j_{14}+j_{15})}$ . However, for the  $10j$  symbol to be non-zero, the spins at each vertex must sum to an integer [4]. Therefore, if the  $10j$  symbol is non-zero, the sign of the integrand will remain positive. The same argument applies if any subset of the  $h_k$  are reversed.

In the integral in equation (3), fully degenerate points occur in 3-dimensional submanifolds of the domain, but we can apply the same gauge-fixing principles as we used to analyse the stationary points, in this case simply fixing  $h_1 = (1, 0, 0, 0)$ . There are then 16 discrete fully degenerate points in the restricted domain: those where  $h_k = \pm h_1$  for  $k = 2, \dots, 5$ . Restricting the integral to the vicinity of these points we obtain

$$(17) \quad \left| \begin{array}{c} \text{Diagram} \\ \text{deg} \end{array} \right| := 16 \int_U \prod_{k<l} K_{\lambda a_{kl}}^R(\phi_{kl}) \frac{dh_2}{2\pi^2} \dots \frac{dh_5}{2\pi^2},$$

where  $U$  is a small open ball around the point  $(h_1, h_1, h_1, h_1) \in (S^3)^4$ .

We approximate this quantity by noting that when  $\phi$  is small, the kernel in (16) is close to:

$$(18) \quad K_a^D(\phi) := \frac{\sin a\phi}{\phi}$$

We call this quantity the ‘degenerate kernel’. As we shall see in Section 5, this is the analog of the original kernel in a spin-network formalism where the space of constant curvature,  $S^3$ , is replaced by three-dimensional Euclidean space, and the group  $\text{SO}(4)$ , the isometry group of  $S^3$ , is replaced by the Euclidean group of isometries of  $\mathbb{R}^3$ .

Similarly, the integral over a small subset of  $(S^3)^4$  can be approximated by an integral over a subset of  $(\mathbb{R}^3)^4$ , and the angle  $\phi_{kl}$  between unit vectors  $h_k$  and  $h_l$  in  $S^3$  replaced by the Euclidean distance  $r_{kl} = |x_k - x_l|$  between vectors  $x_k$  and  $x_l$  in  $\mathbb{R}^3$ . In terms of these new Euclidean variables, the restriction  $h_1 = (1, 0, 0, 0)$  is replaced by  $x_1 = (0, 0, 0)$ .

Thus, when  $\lambda$  is large, we have:

$$(19) \quad \left| \begin{array}{c} \text{Diagram} \\ \text{deg} \end{array} \right| \approx 16 \int_U \prod_{k<l} K_{\lambda a_{kl}}^D(|x_k - x_l|) \frac{dx_2}{2\pi^2} \cdots \frac{dx_5}{2\pi^2},$$

where now  $U$  is a small open ball around the origin of  $(\mathbb{R}^3)^4$ .

The approximation (19) exhibits very simple scaling behaviour. First, note that the degenerate kernel (18) obeys the identity:

$$(20) \quad \begin{aligned} K_{\lambda a}^D(r) &= \frac{\sin \lambda ar}{r} \\ &= \lambda K_a^D(\lambda r) \end{aligned}$$

This allows the scaling of (19) to be deduced from a linear change of variables,  $y_k = \lambda x_k$ :

$$(21) \quad \begin{aligned} &16 \int_U \prod_{k<l} K_{\lambda a_{kl}}^D(|x_k - x_l|) \frac{dx_2}{2\pi^2} \cdots \frac{dx_5}{2\pi^2} \\ &= 16\lambda^{-2} \int_{\lambda U} \prod_{k<l} K_{a_{kl}}^D(|y_k - y_l|) \frac{dy_2}{2\pi^2} \cdots \frac{dy_5}{2\pi^2}, \end{aligned}$$

where  $\lambda U$  is the result of rescaling the neighborhood  $U$  by a factor of  $\lambda$ . If the integral on the right hand side of this equation converges in the limit  $\lambda \rightarrow \infty$ , we obtain the asymptotic formula:

$$(22) \quad \left| \begin{array}{c} \text{Diagram} \\ \text{deg} \end{array} \right| \sim 16\lambda^{-2} \int_{(\mathbb{R}^3)^4} \prod_{k<l} K_{a_{kl}}^D(|y_k - y_l|) \frac{dy_2}{2\pi^2} \cdots \frac{dy_5}{2\pi^2}.$$

We call this integral the ‘degenerate  $10j$  symbol’. Apart from the factor of  $16\lambda^{-2}$ , the integral here can be interpreted as the evaluation of a spin network with edges labelled by unitary irreducible representations of the Euclidean group: the group of isometries of Euclidean 3-space. We discuss this interpretation in more detail in Section 5.

While elegant, the integral in (22) is difficult to compute numerically, very much like the integral for the Lorentzian  $10j$  symbol. To obtain a more convenient form for evaluation, we make use of the Kirillov trace formula:

$$\frac{\sin |x|}{|x|} = \frac{1}{4\pi} \int_{S^2} \exp(ix \cdot \xi) d\xi,$$

where  $x$  is a vector in  $\mathbb{R}^3$ , and the integral is over the unit 2-sphere with its standard measure. This formula is easily confirmed by choosing spherical coordinates such that the  $z$ -axis is



parallel to the vector  $x$ . For our purposes we shall rewrite it as follows:

$$(23) \quad K_a^D(|x|) = \int_{S(a)} \exp(ix \cdot \xi) d\xi,$$

where  $S(a)$  is the 2-sphere of radius  $a$  embedded in  $\mathbb{R}^3$ , but where  $d\xi$  is the induced Lebesgue measure divided by  $4\pi a$ , to hide some annoying constants that would otherwise appear in this formula. Using this we can rewrite (22) as:

$$(24) \quad \left| \begin{array}{c} \text{graph} \\ \text{deg} \end{array} \right| \sim \frac{16\lambda^{-2}}{(2\pi^2)^4} \int_{(\mathbb{R}^3)^4} \int_{\mathcal{X}} \exp(i \sum_{k<l} (y_k - y_l) \cdot \xi_{kl}) d\xi_{12} \cdots d\xi_{45} dy_2 \cdots dy_5$$

$$= \frac{\lambda^{-2}}{\pi^8} \int_{(\mathbb{R}^3)^4} \int_{\mathcal{X}} \exp(i(y_2, y_3, y_4, y_5) \cdot F(\xi_{12}, \dots, \xi_{45})) d\xi_{12} \cdots d\xi_{45} dy_2 \cdots dy_5$$

$$= \lambda^{-2} 2^{12} \pi^4 \int_{\mathcal{X}} \delta^{12}(F(\xi_{12}, \dots, \xi_{45})) d\xi_{12} \cdots d\xi_{45},$$

where

$$\mathcal{X} = \prod_{k<l} S(a_{kl})$$

is a Cartesian product of 2-spheres with the measures described above, and  $F: \mathcal{X} \rightarrow \mathbb{R}^{12}$  is defined by:

$$F(\xi_{12}, \dots, \xi_{45}) = \begin{aligned} &(-\xi_{12} + \xi_{23} + \xi_{24} + \xi_{25}, \\ &-\xi_{13} - \xi_{23} + \xi_{34} + \xi_{35}, \\ &-\xi_{14} - \xi_{24} - \xi_{34} + \xi_{45}, \\ &-\xi_{15} - \xi_{25} - \xi_{35} - \xi_{45}). \end{aligned}$$

If

$$\mathcal{N} = \{\xi \in \mathcal{X} : F(\xi) = 0\}$$

then the final integral in (24) will be well-behaved so long as at each point  $\xi \in \mathcal{N}$  the differential of  $F$  has maximal rank, namely 12. When this is the case  $\mathcal{N}$  is an 8-dimensional submanifold of  $\mathcal{X}$ , and the integral reduces to:

$$(25) \quad \left| \begin{array}{c} \text{graph} \\ \text{deg} \end{array} \right| \sim \lambda^{-2} 2^{12} \pi^4 \int_{\mathcal{N}} |J(\xi)|^{-1} d\xi$$

Here  $d\xi$  is the Lebesgue measure on  $\mathcal{N}$  induced by the Riemannian metric on  $\mathcal{X}$ , but divided by a factor of  $\prod_{k<l} 4\pi a_{kl}$ , since we have divided the Lebesgue measure on each sphere by a factor of  $4\pi a_{kl}$ . If we choose local coordinates  $(x^1, \dots, x^{20})$  on  $\mathcal{X}$  near  $\xi \in \mathcal{N}$  such that  $x^1, \dots, x^8$  are zero on  $\mathcal{N}$ , then  $|J(\xi)|$  is the Jacobian determinant of  $F$  as a function of the remaining 12 variables  $x^9, \dots, x^{20}$ .

One way to get solutions of  $F = 0$  is to start with 4-simplices in  $\mathbb{R}^3$ : that is, 5-tuples of points in  $\mathbb{R}^3$ , together with the ten triangles and five tetrahedra determined by these points. Given such a 4-simplex, let  $\xi_{kl}$  be the vector that is normal to the triangle shared by the  $k$ th and  $l$ th tetrahedra, and has length equal to the area of this triangle. As shown in [16], the four vectors  $\xi_{kl}$  normal to the triangles in any one tetrahedron must sum to zero, with

appropriate signs:

$$(26) \quad \begin{aligned} -\xi_{12} + \xi_{23} + \xi_{24} + \xi_{25} &= 0 \\ -\xi_{13} - \xi_{23} + \xi_{34} + \xi_{35} &= 0 \\ -\xi_{14} - \xi_{24} - \xi_{34} + \xi_{45} &= 0 \\ -\xi_{15} - \xi_{25} - \xi_{35} - \xi_{45} &= 0 \\ \xi_{12} + \xi_{13} + \xi_{14} + \xi_{15} &= 0. \end{aligned}$$

Each vector appears twice in these formulas, with opposite signs, since the outwards-pointing normal to one tetrahedron is the inwards-pointing normal to another. The first four equations say that  $F = 0$ ; the last is an algebraic consequence of the rest. If  $|\xi_{kl}| = a_{kl}$ , the vectors  $\xi_{kl}$  thus determine a point in  $\mathcal{N}$ .

This suggests that we think of points of  $\mathcal{N}$  as ‘degenerate 4-simplices’. However, not every point of  $\mathcal{N}$  comes from a 5-tuple of points in  $\mathbb{R}^3$  this way. To see this, note that two 5-tuples in  $\mathbb{R}^3$  determine the same point in  $\mathcal{N}$  if they are translates of each other. The space of 5-tuples modulo translation has dimension  $3 \times 5 - 3 = 12$ , but the space of solutions of equation (26) has dimension  $30 - 4 \times 3 = 18$ . Thus there are simply not enough 5-tuples of points in  $\mathbb{R}^3$  to account for all solutions of equation (26). We shall still call points of  $\mathcal{N}$  ‘degenerate 4-simplices’, because it is a useful heuristic. However, it is important to take this phrase with a grain of salt.

Of course, for the space  $\mathcal{N}$  to be nonempty, there must *exist* a solution of  $F = 0$ . This imposes certain restrictions on the numbers  $a_{kl}$ . For example, there will be a solution of  $\xi_{12} + \xi_{13} + \xi_{14} + \xi_{15} = 0$  if and only if these ‘tetrahedron inequalities’ hold:

$$\begin{aligned} a_{12} &\leq a_{13} + a_{14} + a_{15} \\ a_{13} &\leq a_{12} + a_{14} + a_{15} \\ a_{14} &\leq a_{12} + a_{13} + a_{15} \\ a_{15} &\leq a_{12} + a_{13} + a_{14}. \end{aligned}$$

In general  $\mathcal{N}$  will be empty if the four numbers  $a_{kl}$  corresponding to the faces of any one tetrahedron violate the tetrahedron inequalities. In this case the degenerate  $10j$  symbols vanish. Similarly, the Riemannian  $10j$  symbols vanish if the spins  $j_{kl}$  violate the tetrahedron inequalities [4], and for the same sort of geometrical reason [16].

For computational purposes it is helpful to rewrite the integral for the degenerate  $10j$  symbols in yet another way. To do this, first we note that we can exploit  $\text{SO}(3)$  invariance to convert (25) to an integral over the 5-dimensional manifold  $\mathcal{N}/\text{SO}(3)$ , by rotating the  $\xi_{kl}$  in  $\mathcal{X}$  so that they lie in the 17-dimensional subspace where  $\xi_{23}$  is fixed at  $(0, 0, a_{23})$ , and  $\xi_{34}$  lies in the  $x > 0$  half of the  $xz$ -plane. This corresponds to integrating over all possible values for  $\xi_{23}$ , inserting a factor equal to the volume of  $S(a_{23})$  with our chosen measure, which is  $a_{23}$ , and also integrating over all possible azimuthal coordinates for  $\xi_{34}$  and inserting a factor of  $2\pi$ .

In what follows, we will choose coordinates so that the  $12 \times 12$  matrix  $J(\xi)$  whose determinant we require is block diagonal, with two  $3 \times 3$  blocks and one  $6 \times 6$  block.

It turns out to be convenient to parameterise  $\xi_{34}$ , not by its angle from the  $z$ -axis, but by the length

$$s_1 := |\xi_{23} - \xi_{34}|.$$

Since  $\xi_{13} + \xi_{23} = \xi_{34} + \xi_{35}$ , these four vectors can be positioned to form the sides of a (possibly non-planar) quadrilateral. Then  $s_1$  is the length of one of the diagonals. With the vectors  $\xi_{23}$  and  $\xi_{34}$  fixed and the lengths of  $\xi_{13}$  and  $\xi_{35}$  specified, the only remaining freedom this quadrilateral has, if it is to remain closed, is the ‘hinge angle’,  $\alpha_1$ , between the two triangles that meet along the diagonal. The vectors  $\xi_{13}$  and  $\xi_{35}$  have 4 degrees of freedom in all, and

specifying  $\alpha_1$  removes one of them, leaving 3 which break the quadrilateral. With our chosen measure on  $\mathcal{N}$ , the product of the measure for the coordinates we are integrating over, and the Jacobian determinant for the 3 that break the quadrilateral, is:

$$\frac{1}{(4\pi)^3 a_{23}}.$$

We can treat a second 4-tuple of vectors more or less identically. If the vector  $\xi_{12}$  has an azimuthal angle of  $\phi$ , and we define

$$s_2 := |\xi_{23} - \xi_{12}|,$$

then the quadrilateral formed by  $\xi_{23}, \xi_{12}, \xi_{25}$  and  $\xi_{24}$  can be assigned a ‘hinge angle’ of  $\alpha_2$ . This specifies all the degrees of freedom that allow this quadrilateral to remain closed. Once again, the product of the measure for the coordinates we are integrating over, and the Jacobian determinant for the 3 that break the quadrilateral, is:

$$\frac{1}{(4\pi)^3 a_{23}}.$$

So far, we have specified 5 degrees of freedom for  $\mathcal{N}/\text{SO}(3)$ :  $s_1, s_2, \alpha_1, \alpha_2$  and  $\phi$ . No further continuous degrees of freedom remain. The three vectors we have yet to specify must form triangles that complete two quadrilaterals with vectors that have already been parameterised, and because these two triangles have a vector in common, there is no ‘hinge’ freedom left.

Specifically, the three vectors  $\xi_{14}, \xi_{15}$  and  $\xi_{45}$  must form a tetrahedron by fitting over a triangular base that has, as two of its sides, the vectors:

$$\begin{aligned} v &:= \xi_{35} + \xi_{25} \\ w &:= -\xi_{24} - \xi_{34}. \end{aligned}$$

Given their fixed lengths, this determines  $\xi_{14}, \xi_{15}$  and  $\xi_{45}$  completely, apart from the freedom to locate the apex of the tetrahedron on either side of the plane spanned by  $v$  and  $w$ . This freedom can be accounted for with a factor of 2 in the integral. The final contribution to the Jacobian comes from a  $6 \times 6$  block involving all the coordinates of  $\xi_{14}, \xi_{15}$  and  $\xi_{45}$ , and with our chosen measure this is:

$$\frac{1}{(4\pi)^3 6V(a_{14}, a_{15}, a_{45}, |w|, |v-w|, |v|)},$$

where  $V$  is the volume of the tetrahedron as a function of its edge lengths.

Combining these results, we can rewrite (25) as:

$$(27) \quad \left| \begin{array}{c} \text{tetrahedron} \\ \text{deg} \end{array} \right| \sim \frac{\lambda^{-2}}{96\pi^4 a_{23}} \int_{s_1} \int_{s_2} \int_{\alpha_1} \int_{\alpha_2} \int_{\phi} \frac{ds_1 ds_2 d\alpha_1 d\alpha_2 d\phi}{V(a_{14}, a_{15}, a_{45}, |w|, |v-w|, |v|)}.$$

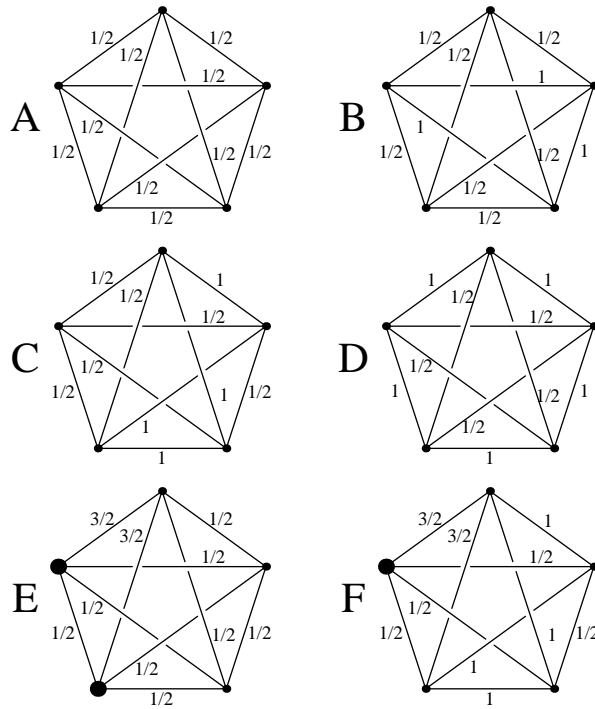
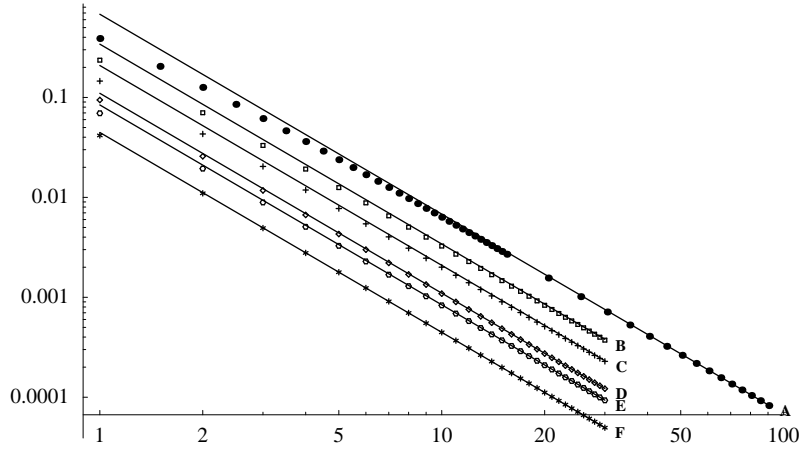
The integrals over the  $s_i$  are taken over intervals determined by the four sides of the quadrilaterals for which they are the diagonal lengths, and the angular variables range from 0 to  $2\pi$ , with the proviso that any part of the domain where the tetrahedron is not geometrically possible must be excluded. In numerical calculations, this can be dealt with by setting the integrand to zero wherever Cayley’s determinant formula for the squared volume of the tetrahedron yields a negative value.

We note that the integrand here is unbounded, and we have not proved that (27) converges, but our numerical calculations suggest that it does.

## 4. NUMERICAL DATA

To test our hypothesis that the asymptotics of  $10j$  symbols are dominated by the contribution of degenerate 4-simplices, we used the algorithm described in [9] to calculate values for several sets of  $10j$  symbols. The figure below shows log-log plots for the absolute values of the Riemannian  $10j$  symbols as a function of  $\lambda$ , where  $\lambda$  is the parameter by which the areas  $a_{kl}$  were multiplied. The legend shows the base spins  $j_{kl}$ ; multiplying  $a_{kl}$  by  $\lambda$  was achieved by replacing the  $j_{kl}$  with:

$$J_{kl} = \lambda j_{kl} + \frac{\lambda - 1}{2}.$$



The lines on the plot show the  $\lambda^{-2}$  asymptotic behaviour described by equation (27); the integrals were evaluated numerically with Lepage's VEGAS algorithm [17], and found to have values of 0.680, 0.341, 0.209, 0.110, 0.0841 and 0.0446 respectively.

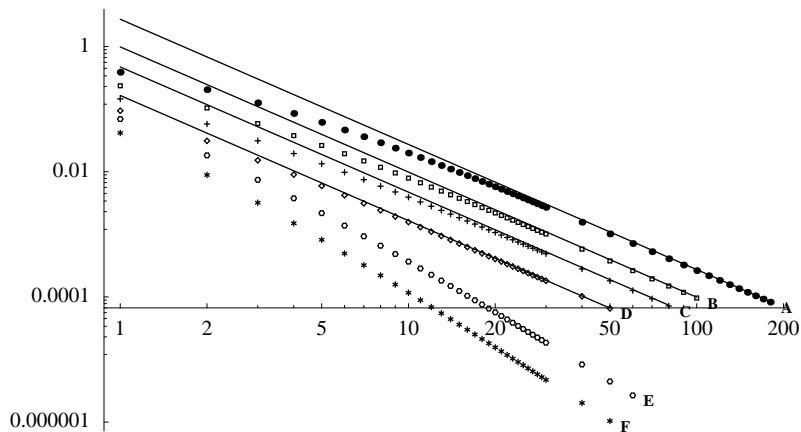
In summary, our numerical data supports the following conjecture. Let us say that the spins  $j_{kl}$  are 'admissible' if for each vertex in the  $10j$  symbol, the spins labelling the four incident edges satisfy the tetrahedron inequalities and sum to an integer. Then:

**Conjecture 1.** *If the ten spins  $j_{kl}$  are admissible, the  $\lambda \rightarrow \infty$  asymptotics of the Riemannian  $10j$  symbols are given by:*

$$\left| \begin{array}{c} \text{Diagram of a 10j symbol (a graph with 10 vertices and 15 edges)} \end{array} \right| \sim 16\lambda^{-2} \int_{(\mathbb{R}^3)^4} \prod_{k<l} K_{2j_{kl}+1}^D(|y_k - y_l|) \frac{dy_2}{2\pi^2} \cdots \frac{dy_5}{2\pi^2}.$$

In the next two sections we generalize this conjecture to a large class of Riemannian and Lorentzian spin networks, including the Lorentzian  $10j$  symbols.

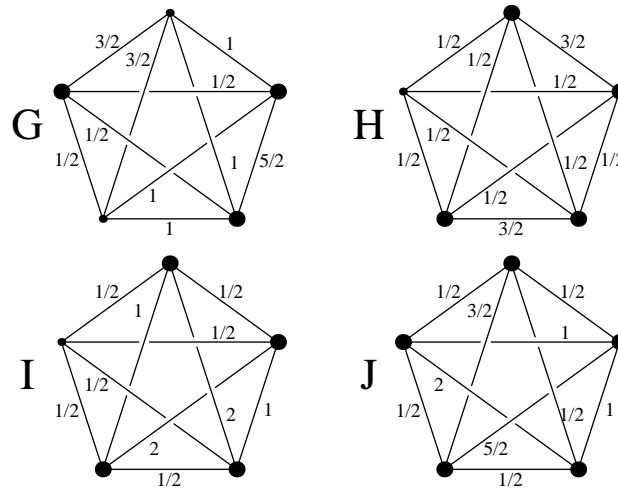
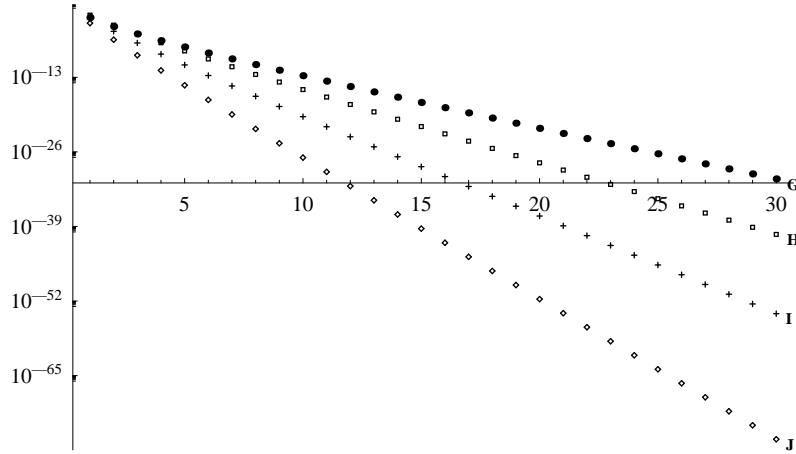
The reader may have wondered why we consider asymptotics of the  $10j$  symbols as *areas* are rescaled, instead of *spins*. The reason is that they are much simpler. We also calculated values for sets of  $10j$  symbols where the spins  $j_{kl}$  were multiplied by  $\lambda$ . The figure below shows log-log plots for the absolute values of the Riemannian  $10j$  symbols with spins  $\lambda j_{kl}$ , where the  $j_{kl}$  match those shown in the legend for the previous figure:



Since the areas  $2\lambda j_{kl} + 1$  are asymptotic to a series of values  $2\lambda j_{kl}$  that are proportional to  $\lambda$ , the data might be expected to exhibit  $\lambda^{-2}$  scaling. And in fact, in cases A-D, that is exactly what is seen. Furthermore, numerically evaluating the integral (27) with  $a_{kl} = 2j_{kl}$  provided the correct coefficients for lines on the plot which are asymptotic to the data; these coefficients were 2.73, 0.987, 0.472 and 0.165 respectively.

For cases E and F, the  $10j$  symbols have one or more vertices at the 'border of admissibility': that is, vertices where three of the spins labelling incident edges sum to equal the fourth spin, making one of the tetrahedron inequalities a strict equality. These vertices are marked with heavy dots on the legend. Setting  $a_{kl} = 2j_{kl}$  in (27) in these cases yields a coefficient of zero, because the domain for at least one variable of integration, the quadrilateral diagonal length  $s_1$ , was reduced to a single point. Empirically, the data here appears to scale as  $\lambda^{-3}$ .

While  $10j$  symbols with one or two vertices on the border of admissibility have  $\lambda^{-3}$  asymptotics with spin rescaling,  $10j$  symbols with three or more vertices on the border of admissibility appear to decay exponentially as their spins are multiplied by  $\lambda$ , as illustrated in the log-linear plot below:



To prove convergence of the partition function in our newly formulated version of the Barrett–Crane model [6], it will probably be necessary to understand the  $\lambda^{-3}$  and exponential decay of borderline-admissible  $10j$  symbols under spin rescaling. However, these are more delicate phenomena than we are prepared to tackle here.

### 5. DEGENERATE SPIN NETWORKS

Though the Riemannian  $10j$  symbols are given by an integral over a product of copies of  $S^3$ , we have seen that when we rescale the areas by a large constant  $\lambda$ , this integral is dominated by the contribution of a very small patch of this space, which can be approximated by a product of copies of  $\mathbb{R}^3$ . Indeed, by a change of variables we can think of the  $\lambda \rightarrow \infty$

limit as one in which the radius of  $S^3$  approaches infinity, so that it degenerates to Euclidean 3-space.

The Riemannian  $10j$  symbols can also be described using the representation theory of  $\text{SO}(4)$ , the isometry group of  $S^3$ . This suggests that the  $\lambda \rightarrow \infty$  asymptotics of the  $10j$  symbols can be described in terms of the representation theory of the Euclidean group  $\text{E}(3)$ , the isometry group of  $\mathbb{R}^3$ .

To do this, we introduce certain spin networks associated to  $\text{E}(3)$  which we call ‘degenerate spin networks’. These are more closely analogous to the Lorentzian spin networks defined in [12] than to the Riemannian spin networks we have been discussing so far, because the edge labels are not restricted to discrete values, and the group representations are infinite-dimensional. However, all three sorts of spin network form part of a unified theory, as outlined below:

geometry	signature	symmetry group	homogeneous space
Riemannian	(++++)	$\text{SO}(4)$	$S^3$
degenerate	(0+++)	$\text{E}(3)$	$\mathbb{R}^3$
Lorentzian	(-+++)	$\text{SO}(3, 1)$	$H^3$

The rotation group  $\text{SO}(4)$  and the Lorentz group  $\text{SO}(3, 1)$  consist of linear transformations of  $\mathbb{R}^4$  with determinant 1 that preserve metrics of signature (++++) and (-+++), respectively. Similarly, the Euclidean group  $\text{E}(3)$  is isomorphic to the group of linear transformations of  $\mathbb{R}^4$  with determinant 1 that preserve the singly degenerate metric with signature (0+++). The groups  $\text{SO}(4)$  and  $\text{SO}(3, 1)$  both ‘contract’ to  $\text{E}(3)$ , meaning that they have it as a limiting case if one forms the isometry group of the metric  $\text{diag}(\epsilon, 1, 1, 1)$  and lets  $\epsilon \downarrow 0$  and  $\epsilon \uparrow 0$ , respectively.

This makes it plausible that the asymptotics of not only Riemannian but also *Lorentzian* spin networks can be calculated using degenerate spin networks. In this section we state a precise conjecture along these lines for a large class of Riemannian spin networks, and present some supporting evidence. In Section 6 we do the same for Lorentzian spin networks.

We begin by describing degenerate spin networks and how to evaluate them. The representations  $j \otimes j$  labelling edges of a Riemannian spin network are representations not just of  $\text{Spin}(4)$ , but actually of  $\text{SO}(4)$ . As emphasized by Freidel and Krasnov [13], these representations are precisely the eigenspaces of the Laplacian on  $S^3$ , which is the homogeneous space  $\text{SO}(4)/\text{SO}(3)$ . Similarly, the representations labelling edges of a Lorentzian spin network are the eigenspaces of the Laplacian on hyperbolic 3-space,  $H^3 = \text{SO}_0(3, 1)/\text{SO}(3)$  — here we must take the connected component of the Lorentz group to get just one sheet of the hyperboloid. Following this pattern, the representations labelling edges of a degenerate spin network should be the eigenspaces of the Laplacian on  $\mathbb{R}^3 = \text{E}(3)/\text{SO}(3)$ .

In fact, there is one representation of this sort for each positive real number  $a$ . Any complex function on  $\mathbb{R}^3$  with  $\nabla^2 f = -a^2 f$  can be written as:

$$(28) \quad f(x) = \int_{\xi \in S(a)} \hat{f}(\xi) \exp(i\xi \cdot x) d\xi,$$

where  $S(a)$  is the 2-sphere of radius  $a$  centered at the origin of  $\mathbb{R}^3$ , and  $d\xi$  is the induced Lebesgue measure divided by  $4\pi a$ . Defining an inner product on these solutions by

$$\langle f, g \rangle = \int_{\xi \in S(a)} \hat{f}(\xi) \overline{\hat{g}(\xi)} d\xi,$$

we form a Hilbert space  $\mathcal{H}_a$  consisting of all solutions  $f$  with  $\langle f, f \rangle < \infty$ . This Hilbert space becomes a representation of the Euclidean group where each group element  $g$  acts via  $(gf)(x) = f(g^{-1}x)$ . In fact, this representation is unitary and irreducible. The functions  $\exp(i\xi \cdot x)$  form a ‘basis’, in the sense that any element of  $\mathcal{H}_a$  can be expressed as in equation (28) for a unique square-integrable function  $\hat{f}$  on the sphere.

We define a ‘degenerate spin network’ to be a directed graph with each edge  $e$  labelled by a representation of this form, or equivalently, a positive number  $a_e$ . An intertwiner between tensor products of these representations can be defined at each vertex by taking the product of functions from the representations labelling the *incoming* edges, multiplying it by the product of complex conjugates of functions from representations labelling the *outgoing* edges, and integrating the result over  $\mathbb{R}^3$ . Given this, the standard way to evaluate such a spin network [18] would be to take a ‘trace’: that is, integrate over a basis label  $\xi_e \in S(a_e)$  for each edge of the graph. The result would be:

$$\int_{\prod_{v \in V} \mathbb{R}^3} \prod_{e \in E} \left[ \int_{S(a_e)} \exp(i(x_{s(e)} - x_{t(e)}) \cdot \xi_e) d\xi_e \right] \prod_{v \in V} \frac{dx_v}{2\pi^2}.$$

Here  $E$  denotes the set of edges of the graph,  $V$  denotes its set of vertices, and the vertices  $s(e)$  and  $t(e)$  are the source and target of the edge  $e$ .

However, just as in the Lorentzian case [12], this gives a divergent integral, because the integrand is invariant when we simultaneously translate all the vectors  $x_v \in \mathbb{R}^3$  by the same amount. More generally, if the underlying graph of our spin network consists of several connected components, and we translate the vectors  $x_v$  where  $v$  lies in any one component, the integrand does not change. To keep things simple, let us consider only spin networks whose underlying graph is *connected*. In this case we can sometimes obtain a well-defined integral by ‘gauge-fixing’ one of the vectors rather than integrating over it: that is, setting  $x_{v_1} = x \in \mathbb{R}^3$  for some vertex  $v_1 \in V$ . If we let  $V' = V - \{v_1\}$ , this gives the following formula for evaluating a degenerate spin network with edges labelled by the numbers  $a_e$ :

$$(29) \quad I^D(a) = \int_{\prod_{v \in V'} \mathbb{R}^3} \prod_{e \in E} \left[ \int_{S(a_e)} \exp(i(x_{s(e)} - x_{t(e)}) \cdot \xi_e) d\xi_e \right] \prod_{v \in V'} \frac{dx_v}{2\pi^2}.$$

As in the Lorentzian case [19], one can show that if this integral converges, the result does not depend on our choice of the special vertex  $v_1$  or the point  $x \in \mathbb{R}^3$ . Assuming the integral does converge, we can use the Kirillov trace formula (23) to reexpress it as:

$$(30) \quad I^D(a) = \int_{\prod_{v \in V'} \mathbb{R}^3} \prod_{e \in E} [K_{a_e}^D(|x_{s(e)} - x_{t(e)}|)] \prod_{v \in V'} \frac{dx_v}{2\pi^2}.$$

When the evaluation of a degenerate spin network converges, it always obeys a very simple scaling law as we multiply all the edge labels  $a_e$  by the same constant  $\lambda$ . Using the scaling property of the degenerate kernel noted in equation (20), we find that

$$(31) \quad I^D(\lambda a) = \lambda^{|E|-3(|V|-1)} I^D(a),$$

where  $|E|$  is the number of edges in the underlying graph and  $|V|$  is the number of vertices.

We have already argued that the asymptotics of the Riemannian  $10j$  symbols are governed by the corresponding degenerate spin network. We can generalize this argument as follows. Fix a connected graph. If we label each edge  $e$  by a positive integer  $a_e$  — or equivalently a spin  $j_e$  with  $a_e = 2j_e + 1$  — we obtain a Riemannian spin network, whose evaluation we



define by:

$$(32) \quad I^R(a) = \int_{\prod_{v \in V'} S^3} \prod_{e \in E} K_{a_e}^R(d(x_{s(e)}, x_{t(e)})) \prod_{v \in V'} \frac{dx_v}{2\pi^2}.$$

Here  $d(x, y)$  is the distance between points  $x, y$  in the unit 3-sphere in  $\mathbb{R}^4$  as measured by the induced Riemannian metric. This formula is equivalent to the standard integral formula [20], except that we have omitted the usual signs in order to simplify the relationship to degenerate spin networks. Fixing a small open ball  $U$  around some point  $(x, \dots, x) \in \prod_{v \in V'} S^3$  we define the ‘degenerate contribution’ to this integral to be:

$$(33) \quad I_{\text{deg}}^R(a) = 2^{|V|-1} \int_U \prod_{e \in E} K_{a_e}^R(d(x_{s(e)}, x_{t(e)})) \prod_{v \in V'} \frac{dx_v}{2\pi^2}.$$

Just as we included a factor of 16 in equation (17), here we include a factor of  $2^{|V|-1}$  to take into account the contribution of anti-parallel degenerate points; as before there is no cancellation between these if the spins labelling edges incident to each vertex sum to an integer, as they must for the spin network to have a nonzero value. Using the same nonrigorous argument as in Section 3, we see that as  $\lambda \rightarrow \infty$ ,

$$(34) \quad \begin{aligned} I_{\text{deg}}^R(\lambda a) &\sim 2^{|V|-1} \int_U \prod_{e \in E} K_{\lambda a_e}^D(|x_{s(e)} - x_{t(e)}|) \prod_{v \in V'} \frac{dx_v}{2\pi^2} \\ &= 2^{|V|-1} \lambda^{|E|-3(|V|-1)} \int_{\lambda U} \prod_{e \in E} K_{a_e}^D(|y_{s(e)} - y_{t(e)}|) \prod_{v \in V'} \frac{dy_v}{2\pi^2} \\ &\sim 2^{|V|-1} \lambda^{|E|-3(|V|-1)} I^D(a) \\ &= 2^{|V|-1} I^D(\lambda a), \end{aligned}$$

where  $U$  now denotes an open ball around the origin of  $\prod_{v \in V'} \mathbb{R}^3$ , and we made the change of variables  $y_e = \lambda x_e$ .

In short, this argument suggests that the asymptotics of the degenerate contribution to the value of a Riemannian spin network are proportional to those of the corresponding degenerate spin network:

$$I_{\text{deg}}^R(\lambda a) \sim 2^{|V|-1} I^D(\lambda a),$$

and we know the latter are very simple:

$$I^D(\lambda a) = \lambda^{|E|-3(|V|-1)} I^D(a).$$

This is particularly interesting when we also have

$$I^R(\lambda a) \sim I_{\text{deg}}^R(\lambda a),$$

because then we can compute the asymptotics of a Riemannian spin network by evaluating a degenerate spin network:

$$I^R(\lambda a) \sim 2^{|V|-1} \lambda^{|E|-3(|V|-1)} I^D(a).$$

When can we expect this to occur? Clearly we should at least demand that the degenerate contribution outweigh the contribution of stationary phase points. A simple power-counting argument as in Section 2 suggests that the contribution of stationary phase points is of order:

$$(35) \quad I_{\text{stat}}^R(\lambda a) = \begin{cases} O(\lambda^{-\frac{3}{2}|V|+3}) & |V| > 2 \\ O(\lambda^{-\frac{1}{2}}) & |V| = 2 \\ O(\lambda) & |V| = 1, \end{cases}$$

where the graphs with one or two vertices are different because there is less need for ‘gauge-fixing’. Comparing these asymptotics to those of the degenerate contribution, we can formulate the following:

**Conjecture 2.** *Given a connected graph with more than two vertices and  $|E| > \frac{3}{2}|V|$ , or two vertices and  $|E| > 2$ , or one vertex and  $|E| > 1$ , as  $\lambda \rightarrow \infty$  we have*

$$I^R(\lambda a) \sim 2^{|V|-1} \lambda^{|E|-3(|V|-1)} I^D(a)$$

as long as the integral defining  $I^D(a)$  converges and the spins  $j_e$  labelling the edges incident to each vertex are admissible.

Here we say the spins labelling the edges incident to some vertex are ‘admissible’ if they sum to an integer and each is less than or equal to the sum of the rest. We do not yet have general criteria for when the integrals associated to Euclidean spin networks converge, and as we shall see, the relevant theorems are bound to be a bit different than in the Lorentzian case [19].

The simplest test of this conjecture is the ‘theta network’, with two vertices joined by three edges, labelled by positive integers  $a$ ,  $b$ , and  $c$ . When the corresponding spins are admissible, the Riemannian theta network evaluates to:

$$(36) \quad \left( \begin{array}{c} a \\ \bullet \quad \bullet \\ b \\ \bullet \quad \bullet \\ c \end{array} \right)^R = 1.$$

The degenerate theta network can also be explicitly evaluated; assuming without loss of generality that  $a \leq b \leq c$ :

$$(37) \quad \left( \begin{array}{c} a \\ \bullet \quad \bullet \\ b \\ \bullet \quad \bullet \\ c \end{array} \right)^D = \begin{cases} 0 & c > a + b \\ \frac{1}{4} & c = a + b \\ \frac{1}{2} & c < a + b. \end{cases}$$

Since the Riemannian network’s spins are admissible, the third inequality must hold for the corresponding areas  $a, b, c$ . Thus in this case the conjecture gives an *exact* formula for the Riemannian spin network.

The next simplest case is the ‘ $4j$  symbol’: the spin network with two vertices joined by four edges, labelled by positive integers  $a$ ,  $b$ ,  $c$  and  $d$ . Without loss of generality let us assume  $a \leq b \leq c \leq d$ . As noted in a previous paper in this series [6], the Riemannian  $4j$  symbol counts the dimension of a space of  $SU(2)$  intertwiners. Using this it follows that:

$$(38) \quad \left( \begin{array}{c} a \\ \bullet \quad \bullet \\ b \\ \bullet \quad \bullet \\ c \\ \bullet \quad \bullet \\ d \end{array} \right)^R = \begin{cases} 0 & b + c \leq d - a \\ \frac{1}{2}(a + b + c - d) & d - a \leq b + c < d + a \\ a & d + a \leq b + c. \end{cases}$$

The corresponding degenerate spin network evaluates to:

$$(39) \quad \left( \begin{array}{c} a \\ \bullet \quad \bullet \\ b \\ \bullet \quad \bullet \\ c \\ \bullet \quad \bullet \\ d \end{array} \right)^D = \begin{cases} 0 & b + c \leq d - a \\ \frac{1}{4}(a + b + c - d) & d - a \leq b + c < d + a \\ \frac{1}{2}a & d + a \leq b + c, \end{cases}$$

so the conjecture is again exact.

An interesting check on our hypotheses is the tetrahedral spin network. This has four vertices and  $|E| = \frac{3}{2}|V|$ , so the hypotheses of Conjecture 2 do not apply: we expect the stationary phase contribution to the Riemannian tetrahedral network to be comparable to the degenerate contribution. The degenerate tetrahedral network evaluates to:

$$(40) \quad \left( \begin{array}{c} \bullet \\ \diagdown \quad \diagup \\ a \quad c \\ \diagup \quad \diagdown \\ e \quad f \\ \bullet \quad \bullet \\ b \end{array} \right)^D = \frac{1}{24\pi V(a, b, c, d, e, f)},$$

where  $V(a, b, c, d, e, f)$  is defined as the volume of the tetrahedron dual to the tetrahedral network. Each triangle in this dual tetrahedron corresponds to a vertex of the tetrahedral network, and the three sides of the triangle have lengths equal to the labels on the three edges incident to the network vertex:

$$(41) \quad V(a, b, c, d, e, f) = \text{the volume of } \begin{array}{c} \bullet \\ \diagdown \quad \diagup \\ f \quad e \\ \diagup \quad \diagdown \\ c \quad a \\ \bullet \quad \bullet \\ d \end{array}.$$

On the other hand, the Riemannian tetrahedral network evaluates to the square of the  $SU(2)$  tetrahedral network, the basic building-block of the Ponzano–Regge model. Thanks to the calculation of Ponzano and Regge [1], later made rigorous by Roberts [2], this means that:

$$(42) \quad \left( \begin{array}{c} \bullet \\ \diagdown \quad \diagup \\ a \quad c \\ \diagup \quad \diagdown \\ e \quad f \\ \bullet \quad \bullet \\ b \end{array} \right)^R \sim \frac{\cos^2(S + \frac{\pi}{4})}{12\pi V(\frac{a}{2}, \frac{b}{2}, \frac{c}{2}, \frac{d}{2}, \frac{e}{2}, \frac{f}{2})} \\ = \frac{1 + \cos 2(S + \frac{\pi}{4})}{24\pi V(\frac{a}{2}, \frac{b}{2}, \frac{c}{2}, \frac{d}{2}, \frac{e}{2}, \frac{f}{2})}.$$

Here we are dealing with a dual tetrahedron whose edge lengths are  $\frac{1}{2}a_{kl}$ , where  $a_{kl}$  ranges over  $a, b, c, d, e, f$  as  $1 \leq k < l \leq 4$ . This tetrahedron has Regge action

$$S = \sum_{1 \leq k < l \leq 4} \frac{1}{2} a_{kl} \theta_{kl},$$

where  $\theta_{kl}$  are the corresponding dihedral angles. Thus it appears that the asymptotics of the Riemannian tetrahedral network are a sum of two parts: a part equal to 8 times the degenerate tetrahedral network, and an oscillatory part coming from the stationary phase points.

When the edge lengths of the dual tetrahedron are such that it cannot exist in Euclidean space, the degenerate tetrahedral network evaluates to zero, and the Riemannian network obeys different asymptotics in which its evaluation exponentially decays with increasing spin.

## 6. LORENTZIAN SPIN NETWORKS

Now let us turn to Lorentzian spin networks [12, 19]. Each positive real number determines a unitary irreducible representation of the connected Lorentz group  $\mathrm{SO}_0(3, 1)$  corresponding to an eigenspace of the Laplacian on  $H^3$ ; however, we prefer to describe the evaluation using an integral formula. Fixing a connected graph with vertex set  $V$  and edge set  $E$ , and labelling each edge  $e$  by a positive real number  $a_e$ , we obtain a so-called ‘Lorentzian spin network’, whose evaluation is given by:

$$(43) \quad I^L(a) = \int_{\prod_{v \in V'} H^3} \prod_{e \in E} K_{a_e}^L(d(x_{s(e)}, x_{t(e)})) \prod_{v \in V'} \frac{dx_v}{2\pi^2}.$$

As with degenerate spin networks, we have chosen a vertex  $v_1$ , set  $V' = V - \{v_1\}$ , and let  $x_{v_1}$  be any fixed point in  $H^3$ . Here  $H^3$  is hyperbolic 3-space, i.e., the submanifold of Minkowski spacetime given by:

$$H^3 = \{t^2 - x^2 - y^2 - z^2 = 1, t > 0\}$$

with its induced Riemannian metric. We define the Lorentzian kernel  $K^L$  by:

$$(44) \quad K_a^L(\phi) := \frac{\sin a\phi}{\sinh \phi}.$$

We warn the reader that this convention differs from that of most previous papers [5, 12, 19], which include a factor of  $a$  in the denominator. Including that factor would divide any Lorentzian spin network by the product of its edge labels, so for example, it would divide the asymptotics of the Lorentzian  $10j$  symbols as defined here by a factor of  $\lambda^{10}$ .

The same line of argument by which we arrived at our conjecture concerning asymptotics of Riemannian spin networks applies to Lorentzian ones. The most important difference is that no factor of  $2^{|V|-1}$  appears, since there are no ‘antipodal points’ in hyperbolic space. We have not investigated criteria for the existence of stationary phase points, but where they are present their exponents will be the same as in the Riemannian case, leading us to make:

**Conjecture 3.** *Given a connected graph with more than two vertices and  $|E| > \frac{3}{2}|V|$ , or two vertices and  $|E| > 2$ , or one vertex and  $|E| > 1$ , as  $\lambda \rightarrow \infty$  we have*

$$I^L(\lambda a) \sim \lambda^{|E|-3(|V|-1)} I^D(a)$$

as long as the integral defining  $I^D(a)$  converges and the positive numbers  $a_e$  labelling edges incident to each vertex are admissible.

Here we say the positive numbers labelling the edges incident to some vertex are ‘admissible’ if each is strictly less than the sum of the rest.

Again the simplest test of this conjecture is the Lorentzian theta network. Translating their result into our notation, a calculation of Barrett and Crane [12] shows that for any  $a, b, c > 0$ ,

$$\left( \begin{array}{c} \text{---} a \text{---} \\ \text{---} b \text{---} \\ \text{---} c \text{---} \end{array} \right)^L = \frac{1}{4} [f(-a + b + c) + f(a - b + c) + f(a + b - c) - f(a + b + c)],$$

where

$$f(k) = \tanh\left(\frac{\pi}{2}k\right).$$

As the conjecture predicts, the asymptotics of this match those of the degenerate theta network, which are given by:

$$\left( \begin{array}{c} a \\ \text{---} \\ b \\ \text{---} \\ c \end{array} \right)^D = \frac{1}{4} [\text{sign}(-a + b + c) + \text{sign}(a - b + c) + \text{sign}(a + b - c) - \text{sign}(a + b + c)].$$

It is worth noting that while the integral for the Lorentzian theta network converges even when we take the absolute value of the integrand, this fails for the degenerate theta network. This makes it more challenging to find criteria for convergence of degenerate spin networks, since we cannot simply mimic the theory that applies in the Lorentzian case [19].

Barrett and Crane also worked out the Lorentzian  $4j$  symbols, obtaining:

$$\left( \begin{array}{c} a \\ \text{---} \\ b \\ \text{---} \\ c \\ \text{---} \\ d \end{array} \right)^L = \frac{1}{4} [g(-a + b + c + d) + g(a - b + c + d) + g(a + b - c + d) + g(a + b + c - d) \\ - g(a + b - c - d) - g(a - b + c - d) - g(a - b - c + d) - g(a + b + c + d)],$$

where

$$g(k) = \frac{k}{2} \coth\left(\frac{\pi}{2}k\right).$$

From equation (39) one can show:

$$\left( \begin{array}{c} a \\ \text{---} \\ b \\ \text{---} \\ c \\ \text{---} \\ d \end{array} \right)^D = \frac{1}{4} [h(-a + b + c + d) + h(a - b + c + d) + h(a + b - c + d) + h(a + b + c - d) \\ - h(a + b - c - d) - h(a - b + c - d) - h(a - b - c + d) - h(a + b + c + d)],$$

where

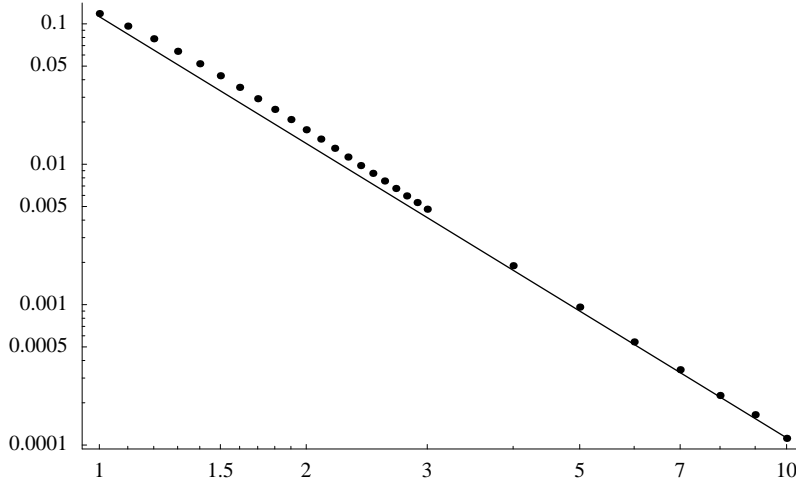
$$h(k) = \frac{k}{2} \text{sign}(k).$$

Thus the conjecture is also confirmed in this case.

Next, consider the tetrahedral spin network. As in the Riemannian case, Conjecture 3 does not apply, but we can still predict the asymptotic behaviour of the contribution of the fully degenerate point:

$$\left( \begin{array}{c} \bullet \\ \lambda \quad \lambda \\ \lambda \quad \lambda \\ \lambda \end{array} \right)^L \sim \frac{1}{24\pi V(\lambda, \lambda, \lambda, \lambda, \lambda, \lambda)} + SP \\ = \frac{\sqrt{2}}{4\pi} \lambda^{-3} + SP,$$

where ‘ $SP$ ’ represents the contribution from stationary phase points, if any. Below is a log-log plot comparing this prediction to numerical data. The horizontal axis in this graph represents  $\lambda$ , while the vertical axis represents the value of the tetrahedral spin network.



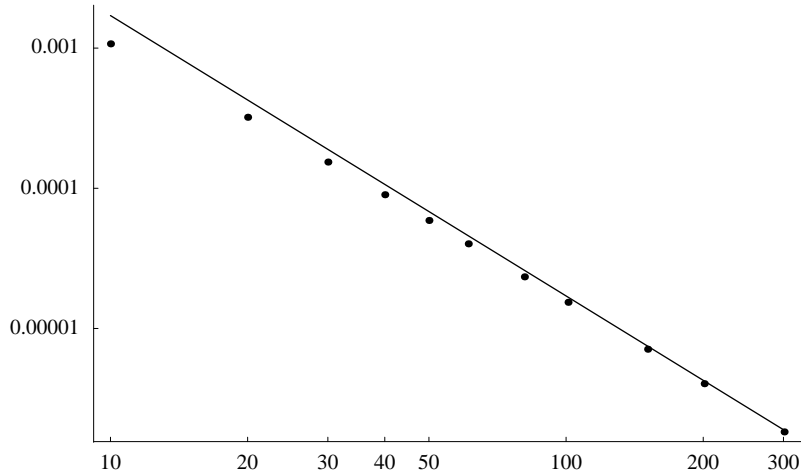
The most interesting test of Conjecture 3 is the  $10j$  symbol. If the conjecture is true, the Lorentzian  $10j$  symbol should be asymptotic to the degenerate  $10j$  symbol, and therefore asymptotic to  $1/16$  times the Riemannian  $10j$  symbol. Since the Riemannian  $10j$  symbol is positive [5], this in turn would imply that the Lorentzian and degenerate  $10j$  symbols are positive in the  $\lambda \rightarrow \infty$  limit.

It is difficult to compute the Lorentzian  $10j$  symbol, but we have numerically checked the conjecture in the special case where all the edges are labelled by the same number  $\lambda$ . In this case the conjecture says that:

$$\left( \begin{array}{c} \lambda \\ \lambda \quad \lambda \\ \lambda \quad \lambda \quad \lambda \\ \lambda \quad \lambda \quad \lambda \\ \lambda \end{array} \right)^L := \int_{(H^3)^4} \prod_{k < l} K_\lambda^L(\phi_{kl}) \frac{dx_2}{2\pi^2} \cdots \frac{dx_5}{2\pi^2}$$

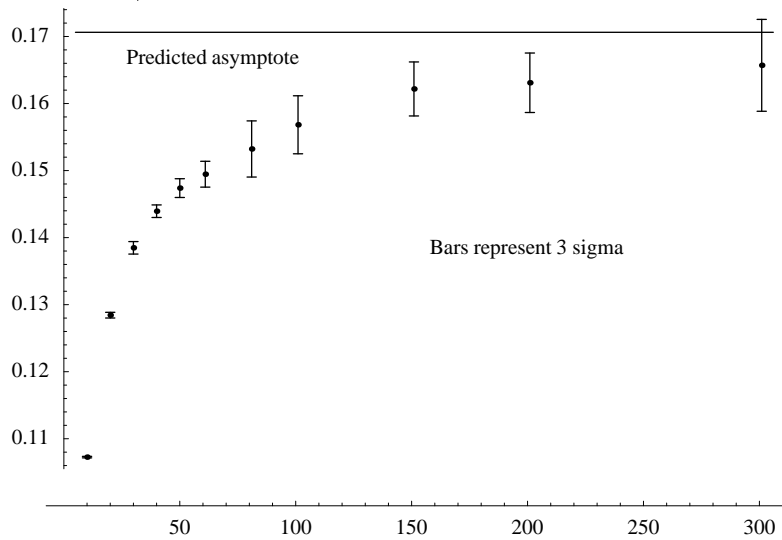
$$\sim .1706 \lambda^{-2},$$

where of course the constant is not exact. Below we show a log-log plot comparing this prediction to data points obtained by computing the Lorentzian  $10j$  symbols numerically. The horizontal axis represents  $\lambda$ , while the vertical axis represents the value of the  $10j$  symbol.



Here we computed the Lorentzian  $10j$  symbols by applying the VEGAS algorithm to evaluate the integral above using polar coordinates on  $H^3$ . The integral was reduced from 12 to 9 dimensions by exploiting the  $SO_0(3,1)$  invariance; the infinite domain was made compact by replacing the radial coordinate for each point,  $r_i$ , with a new variable  $t_i = r_i/(1+r_i)$ ; and the domain was further reduced by exploiting a 24-fold symmetry present in the regular case. The large dimension and oscillatory nature of this integral make these calculations extremely computationally intensive.

The graph below shows the same Lorentzian  $10j$  symbols multiplied by  $\lambda^2$ , plotted on a linear scale to give a clearer picture of the rate of convergence towards the asymptote. The error bars are three times the standard deviation computed by the VEGAS algorithm. The coordinate system we used yielded the lowest standard deviations of several we tried, but there is no reason to believe that the estimates of the integral are drawn from a normal distribution, so this data should be treated with caution.



## 7. CONCLUSIONS

It appears that the asymptotics of both Riemannian and Lorentzian  $10j$  symbols can be computed in terms of degenerate spin networks. For the mathematician, this claim still requires proof. Indeed, there is even an issue of rigor concerning our argument that degenerate spin networks asymptotically describe the contribution of fully degenerate points, since we have not proved that the limit

$$\lim_{\lambda \rightarrow \infty} \int_{\lambda U} \prod_{e \in E} K_{a_e}^D(|y_{s(e)} - y_{t(e)}|) \prod_{v \in V'} \frac{dy_v}{2\pi^2}$$

exists. The ambitious mathematician could also try to prove more general versions of our conjectures, in which all the areas  $a_e$  approach infinity, but not in fixed proportion to one another.

For the physicist, however, a more pressing question is: *what do these results imply for the physics of the Barrett–Crane model?* On the one hand, it is unsettling that the simple asymptotic behavior of the Ponzano–Regge model is not found here. The way degenerate geometries govern the asymptotics of the  $10j$  symbols raises the possibility that in the limit of large spins, the Barrett–Crane model reduces to a theory of degenerate metrics. However, it is important to bear in mind that the physics of the Barrett–Crane model may not be controlled by the large-spin behavior of the  $10j$  symbols — though certainly this affects the convergence of the partition function [6]. Indeed, there are still major open questions about the right way to extract physics from spin foam models, and we hope that our work spurs on research on this subject.

**Note Added in Proof.** After we submitted this article for publication, preprints of two forthcoming papers appeared, one by John Barrett and Chris Steele [21] and another by Laurent Freidel and David Louapre [22], both of which deal with the asymptotics of  $10j$  symbols by somewhat different methods than our own, and both of which obtain results in agreement with our conjectures.

**Acknowledgements.** We thank John Barrett and Chris Steele for helpful conversations, and note that we corrected two errors in an earlier draft of this paper (an excess factor of 2 in equation (15) and a false conjecture about stationary phase points in Lorentzian spin network evaluations) after reading a preprint of their forthcoming paper on the asymptotics of  $10j$  symbols.

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## REFERENCES

- [1] G. Ponzano and T. Regge, Semiclassical limit of Racah coefficients, in *Spectroscopic and Group Theoretical Methods in Physics*, ed. F. Bloch, North–Holland, New York, 1968.
- [2] J. Roberts, Classical  $6j$ -symbols and the tetrahedron, *Geom. Top.* **3** (1999), 21–66. Also available as math-ph/9812013.
- [3] J. Roberts, Asymptotics and  $6j$ -symbols, available as math.QA/0201177.
- [4] J. W. Barrett and L. Crane, Relativistic spin networks and quantum gravity, *Jour. Math. Phys.* **39** (1998), 3296–3302. Also available as gr-qc/9709028.
- [5] J. C. Baez and J. D. Christensen, Positivity of spin foam amplitudes, *Class. Quantum Grav.* **19** (2002), 2291–2306. Corrected version available as gr-qc/0110044.
- [6] J. C. Baez, J. D. Christensen, T. R. Halford and D. C. Tsang, Spin foam models of Riemannian quantum gravity, to appear in *Classical and Quantum Gravity*. Available as gr-qc/0202017.
- [7] J. W. Barrett, The classical evaluation of relativistic spin networks, *Adv. Theor. Math. Phys.* **2** (1998), 593–600. Also available as math.QA/9803063.



- [8] J. W. Barrett and R. M. Williams, The asymptotics of an amplitude for the 4-simplex, *Adv. Theor. Math. Phys.* **3** (1999), 209–215. Also available as gr-qc/9809032.
- [9] J. D. Christensen and G. Egan, An efficient algorithm for the Riemannian  $10j$  symbols, *Class. Quantum Grav.* **19** (2002), 1185–1194. Also available as gr-qc/0110045.
- [10] A. Ashtekar, Quantum mechanics of geometry, available as gr-qc/9901023.
- [11] A. Alekseev, A. P. Polychronakos and M. Smedbäck, On area and entropy of a black hole, available as hep-th/0004036.
- [12] J. W. Barrett and L. Crane, A Lorentzian signature model for quantum general relativity, *Class. Quantum Grav.* **17** (2000), 3101–3118. Also available as gr-qc/9904025.
- [13] L. Freidel and K. Krasnov, Simple spin networks as Feynman graphs, *J. Math. Phys.* **41** (2000), 1681–1690. Also available as hep-th/9903192.
- [14] A. Erdelyi, *Asymptotic Expansions*, Dover Publications, New York, 1954.
- [15] H. Mowaffaq, A vector proof of a theorem of Bang, *Amer. Math. Monthly* **108** (2001), 562–564.
- [16] J. C. Baez and J. W. Barrett, The quantum tetrahedron in dimensions 3 and 4, *Adv. Theor. Math. Phys.* **3** (1999), 815–850. Also available as gr-qc/9903060.
- [17] G. P. Lepage, A new algorithm for adaptive multidimensional integration, *Journal of Computational Physics* **27** (1978), 192–203.
- [18] J. C. Baez, An introduction to spin foam models of  $BF$  theory and quantum gravity, in *Geometry and Quantum Physics*, eds. Helmut Gausterer and Harald Grosse, Springer Lecture Notes in Physics **543**, Berlin, 2000. Also available as gr-qc/9905087.
- [19] J. C. Baez and J. W. Barrett, Integrability for relativistic spin networks, *Class. Quant. Grav.* **18** (2001), 4683–4700. Also available as gr-qc/0101107.
- [20] J. W. Barrett, The classical evaluation of relativistic spin networks, *Adv. Theor. Math. Phys.* **2** (1998), 593–600. Also available as math.QA/9803063.
- [21] J. W. Barrett and C. M. Steele, Asymptotics of relativistic spin networks, available as gr-qc/0209023.
- [22] L. Freidel and D. Louapre, Asymptotics of  $6j$  and  $10j$  symbols, available as hep-th/0209134.

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